

A Polar Approach to Fully Localised Planar Patterns

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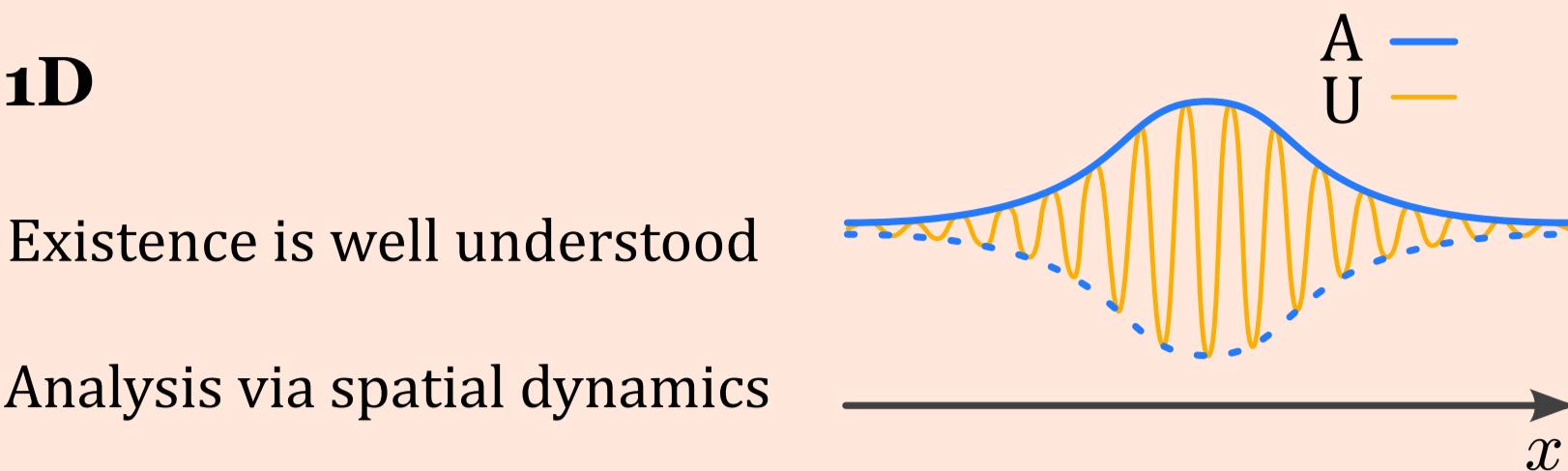


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Planar Localisation



2D

Most physical patterns appear in 2D (or 3D) settings

Localisation in 2 dimensions - no 'time-like' variable for spatial dynamics

Continuum of bifurcating wave numbers $|k| = k$

Localised envelope over periodic domain-covering pattern?



Localised hexagon on the surface of a ferrofluid [LGR+15]

Galerkin Approximation

Truncated expansion

$$u(r, \theta) = \sum_{|n| \leq N} u_{|n|}(r) \cos(mn\theta)$$

Project onto first ($N+1$) modes

Apply radial spatial dynamics to finite dimensional system

2023 - D.J.H, J.J. Bramburger & D.J.B. Lloyd, *Nonlinearity* [HBL23]

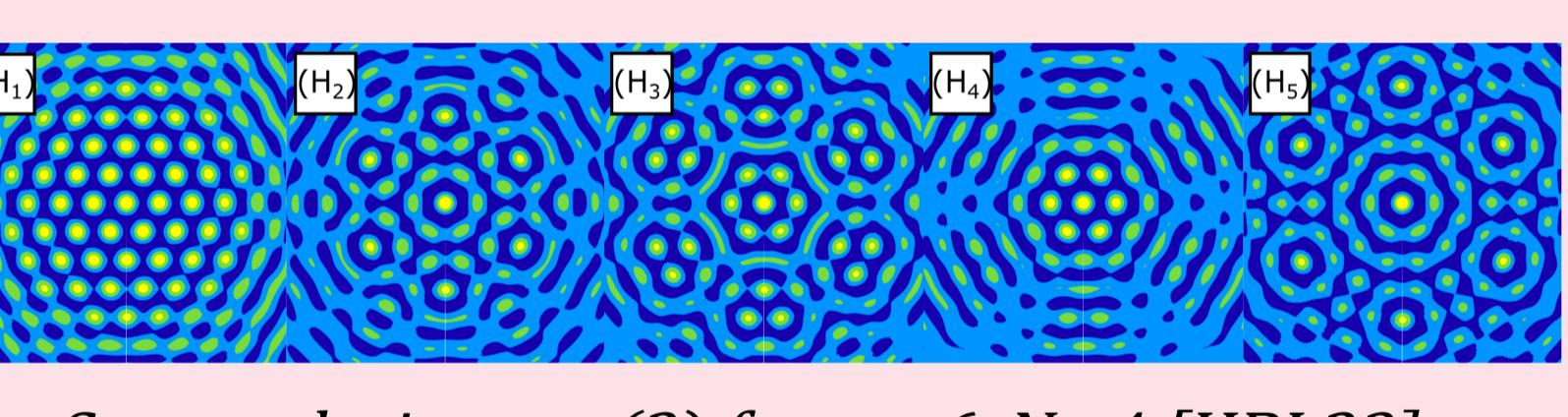
Any solution of the algebraic matching condition

$$a_n = 2 \sum_{j=1}^{N-n} \cos\left(\frac{m\pi(n-j)}{3}\right) a_j a_{n+j} + \sum_{j=0}^n \cos\left(\frac{m\pi(n-2j)}{3}\right) a_j a_{n-j} \quad (2)$$

corresponds to a localised dihedral spot A pattern

2024 - D.J.H, J.J. Bramburger & D.J.B. Lloyd, *Nonlinearity* [HBL24]

Localised dihedral rings with a cubic matching condition



Some solutions to (2) for $m=6, N=4$ [HBL23]

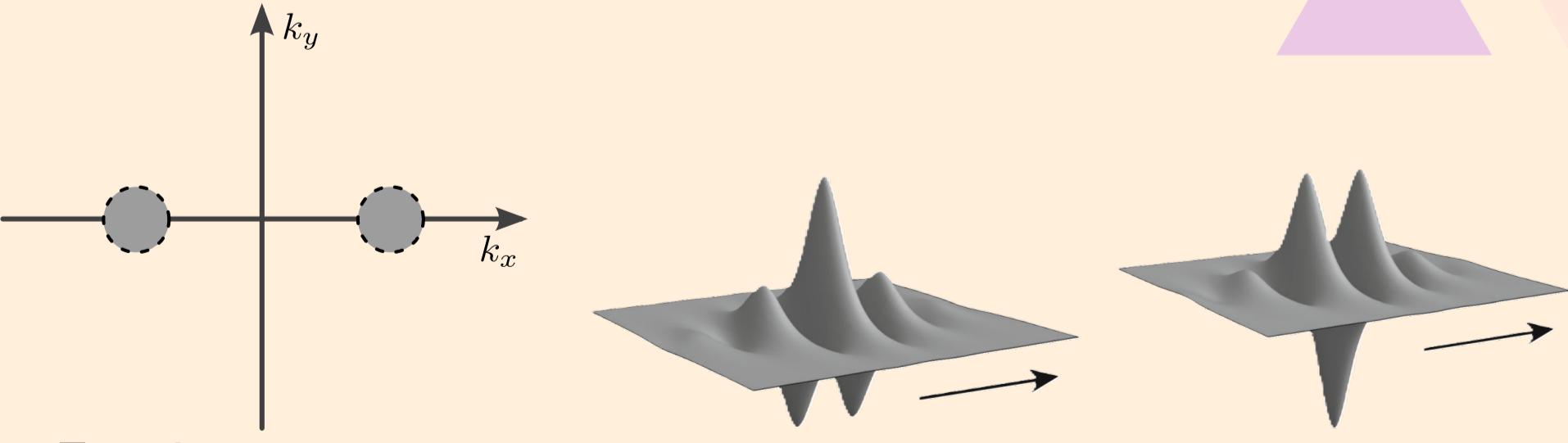
Pushing the Envelope

[Work in progress with M. Groves]

2018 - B. Buffoni et al., *ARMA* [BGW18]

Lyapunov-Schmidt proof of "fully localised" waves

Exponential decay in x and algebraic decay in y



Requires tools from functional analysis and modulation theory

What about fully localised planar patterns in (1)?

2024 - D.J.H. & D.J.B. Lloyd, *SIAM J. Appl. Math.* [HL24]

Formal derivation of radial envelope equations



Replace derivatives with Hankel multipliers

Extend existence proof to radial function spaces

Key Problem

How can we study localised 2D patterns without a spatial dynamics viewpoint?

Can we prove the existence of fully localised patterns in the Swift-Hohenberg equation?

$$\partial_t u = -(1 + \Delta)^2 u - \mu u + \nu u^2 - u^3 \quad (1)$$

Polar Orthogonal Expansions

2D rotation matrix

k -mode function: $\tilde{u}_k(\mathbf{x}) = \frac{1}{2\pi} \int_0^{2\pi} u(\mathbf{R}_k \mathbf{x}) e^{-ik\varphi} d\varphi$

radial k -coefficient: $u_k(r) = \frac{1}{2\pi} \int_0^{2\pi} u(r \cos \varphi, r \sin \varphi) e^{-ik\varphi} d\varphi$

Expansions: $u(\mathbf{x}) = \sum_{n \in \mathbb{Z}} \tilde{u}_n(\mathbf{x}) \quad u(r \cos \theta, r \sin \theta) = \sum_{n \in \mathbb{Z}} u_n(r) e^{in\theta}$

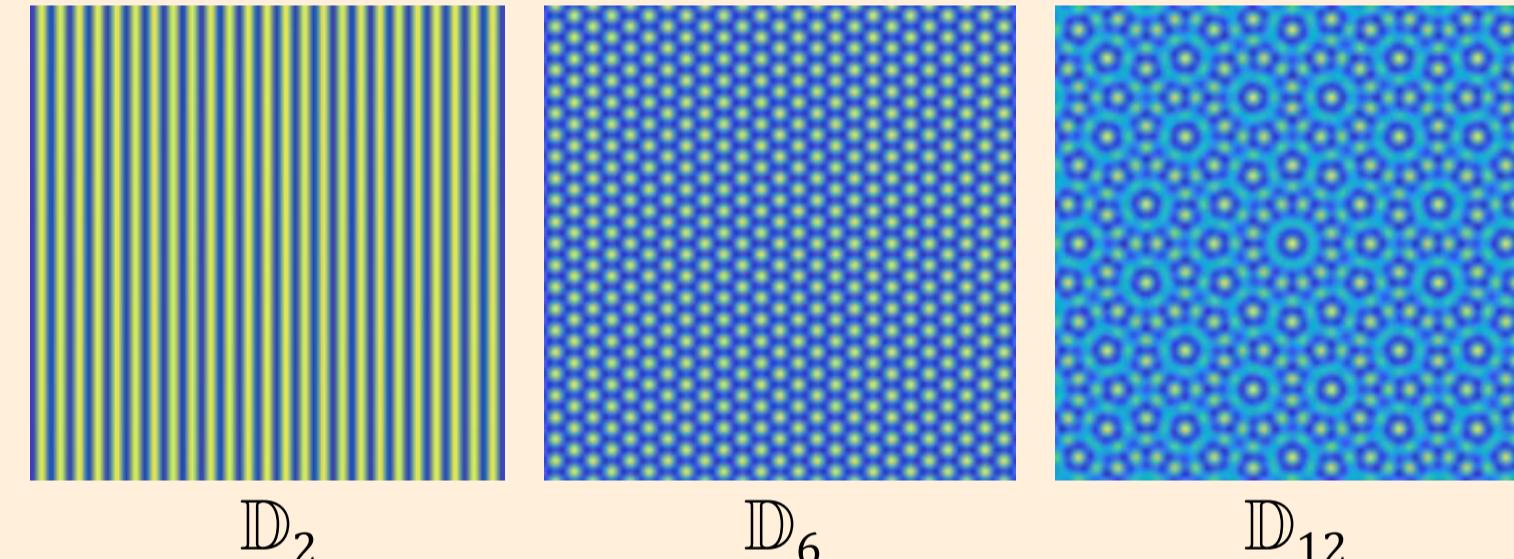
(1) becomes $\left\{ 0 = -(1 + \Delta_n)u_n - \mu u_n + \nu \sum_{i+j=n} u_i u_j - \sum_{i+j+k=n} u_i u_j u_k \right\}_{n \in \mathbb{Z}}$

with n -Laplacian $\Delta_n = (\partial_r + \frac{1-n}{r})(\partial_r + \frac{n}{r})$ and continuity conditions $k u_k(0) = 0$

Reduced expansion for \mathbb{D}_m dihedral functions: $u(r, \theta) = u(r, -\theta) = u(r, \theta + \frac{2\pi}{m})$

Invariant under discrete rotations

$$u(r, \theta) = \sum_{n \in \mathbb{Z}} u_{|n|}(r) \cos(mn\theta)$$



Functions à la mode

2024 - M.D. Groves & D.J.H., *arXiv preprint* [GH24]

k -mode function and radial k -coefficient are related by

$$\tilde{u}_k(r \cos \theta, r \sin \theta) = u_k(r) e^{ik\theta}$$

Wirtinger derivatives: $\partial_\zeta \tilde{u}_k = \frac{1}{\sqrt{2}} (\partial_x - i\partial_y) \quad \partial_{\bar{\zeta}} \tilde{u}_k = \frac{1}{\sqrt{2}} (\partial_x + i\partial_y)$

$$\partial_\zeta \tilde{u}_k = \frac{1}{\sqrt{2}} \mathcal{D}_k u_k e^{i(k-1)\theta} \quad \partial_{\bar{\zeta}} \tilde{u}_k = \frac{1}{\sqrt{2}} \mathcal{D}_{-k} u_k e^{i(k+1)\theta} \quad \mathcal{D}_k = (\partial_r + \frac{k}{r})$$

If u lies in some function space $X(\mathbb{R}^2)$, then \tilde{u}_k lies in the proper subspace $\tilde{X}_{(k)}(\mathbb{R}^2)$ which is isomorphic to the space $X_{(k)}([0, \infty))$ in which u_k lies.

e.g., $u \in C^m(\mathbb{R}^2) \implies \tilde{u}_k \in \tilde{C}_{(k)}^m(\mathbb{R}^2) \iff u_k \in C_{(k)}^m([0, \infty))$

$$u \in C^m(\mathbb{R}^2) \implies \partial_\zeta \tilde{u}_k \in \tilde{C}_{(k-1)}^{m-1}(\mathbb{R}^2) \iff \mathcal{D}_k u_k \in C_{(k-1)}^{m-1}([0, \infty))$$

Fourier transform: $\mathcal{F}[u](\mathbf{k}) = \frac{1}{2\pi} \int_{\mathbb{R}^2} u(\mathbf{x}) e^{-i\mathbf{k} \cdot \mathbf{x}} d\mathbf{x}$

Hankel transform: $\mathcal{H}_k[u_k](\rho) = \int_0^\infty u_k(r) J_k(\rho r) r dr$

$$\mathcal{F}[\tilde{u}_k](\rho \cos \omega, \rho \sin \omega) = i^{-k} \mathcal{H}_k[u_k](\rho) e^{ik\omega}$$

$$H_{(k)}^s([0, \infty)) = \{f \in L_1^2([0, \infty)) : (1 + (\cdot)^2)^{\frac{s}{2}} \mathcal{H}_k[f] \in L_1^2([0, \infty))\}$$

$$L_1^2([0, \infty)) = L^2([0, \infty), r dr)$$

References



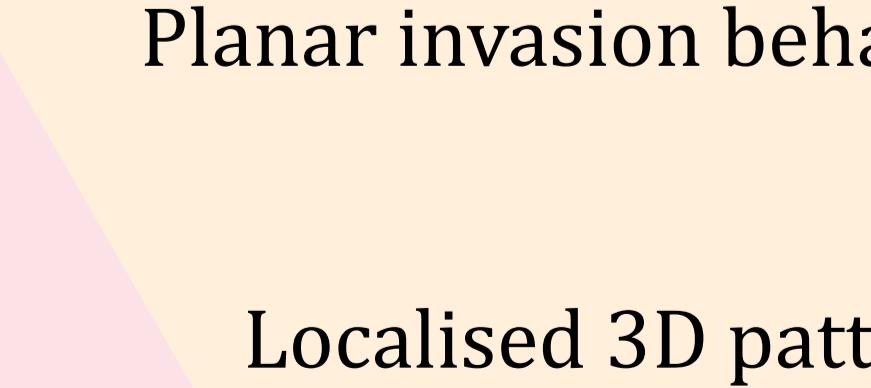
[LGR+15]



[Sch03]



[LS09]



[McS13]



[HBL23]



[HBL24]



[GH24]



[BGW18]



[HL24]

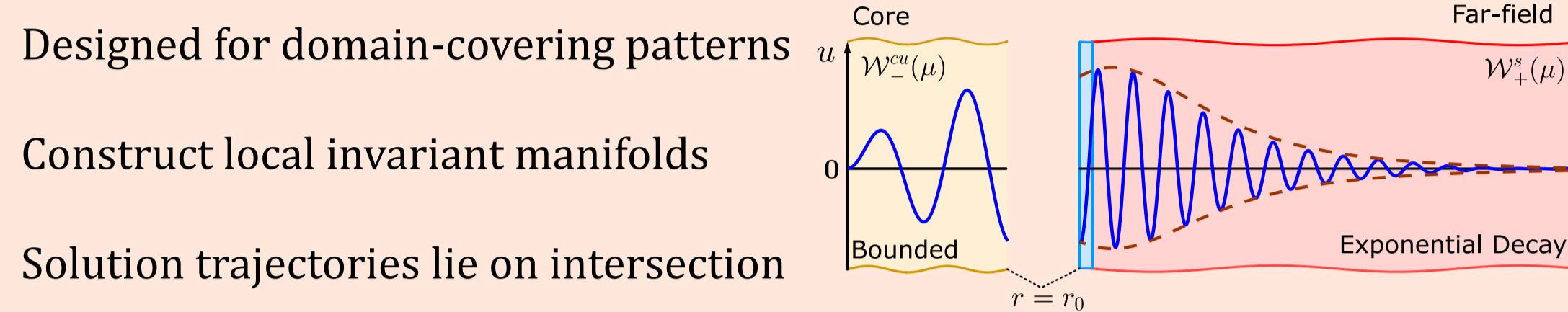
Radial Spatial Dynamics

A simple example: Localised axisymmetric patterns

Invariant under continuous rotations

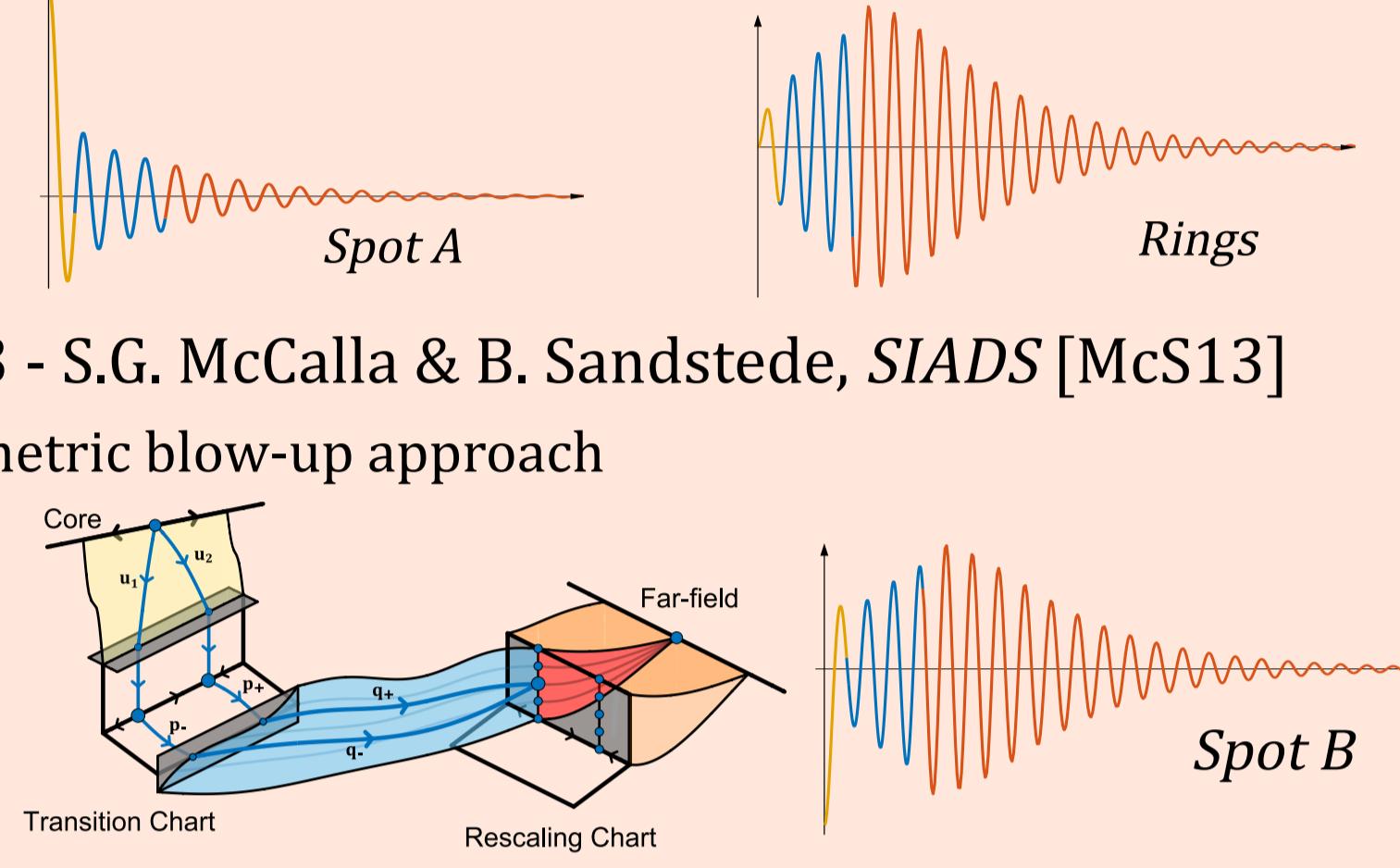
$$u(t, r, \theta) = u_0(r) \text{ - Reduces (1) to a radial ODE}$$

2003 - A. Scheel, *AMS Mem.* [Sch03]



2009 - D.J.B. Lloyd & B. Sandstede, *Nonlinearity* [LS09]

Extension of [Sch03] to localised axisymmetric patterns



Radial Vector Calculus

[Work in progress, upcoming paper]

Variables with physical representation (e.g., fluid velocity) require a choice of basis, which also depends on coordinates

Polar basis: $\mathbf{u}(r, \theta, z) = \sum_{n \in \mathbb{Z}} e^{in\theta} \{u_r^n(r, z) \hat{\mathbf{r}} + u_\theta^n(r, z) \hat{\boldsymbol{\theta}} + u_z^n(r, z) \hat{\mathbf{z}}\}$

$u(r, \theta, z) = \sum_{n \in \mathbb{Z}} e^{in\theta} u^n(r, z) \quad \nabla \mathbf{u} = \sum_{n \in \mathbb{Z}} e^{in\theta} \{\partial_r u^n \hat{\mathbf{r}} + \frac{in}{r} u^n \hat{\boldsymbol{\theta}} + \partial_z u^n \hat{\mathbf{z}}\}$

Complex basis:

$$\hat{\zeta} = \frac{1}{\sqrt{2}} (\hat{\mathbf{x}} + i\hat{\mathbf{y}}) \quad u_\zeta^k = \frac{1}{\sqrt{2}} (u_r^k - iu_\theta^k)$$

$$\hat{\bar{\zeta}} = \frac{1}{\sqrt{2}} (\hat{\mathbf{x}} - i\hat{\mathbf{y}}) \quad u_{\bar{\zeta}}^k = \frac{1}{\sqrt{2}} (u_r^k + iu_\theta^k)$$

$$\mathbf{u} = \sum_{n \in \mathbb{Z}} \{e^{i(n-1)\theta} u_\zeta^n \hat{\zeta} + e^{i(n+1)\theta} u_{\bar{\zeta}}^n \hat{\bar{\zeta}} + e^{in\theta} u_z^n \hat{\mathbf{z}}\}$$

$$\nabla \mathbf{u} = \sum_{n \in \mathbb{Z}} \{e^{i(n-1)\theta} \frac{1}{\sqrt{2}} \mathcal{D}_n u^n \hat{\zeta} + e^{i(n+1)\theta} \frac{1}{\sqrt{2}} \mathcal{D}_{-n} u^n \hat{\bar{\zeta}} + e^{in\theta} \partial_z u^n \hat{\mathbf{z}}\}$$

$$\mathbf{u} \in (X(\mathbb{R}^3))^3 \implies$$

$$(u_\zeta^k, u_{\bar{\zeta}}^k, u_z^k) \in X_{(k-1)}([0, \infty) \times \mathbb{R}) \times X_{(k+1)}([0, \infty) \times \mathbb{R}) \times X_{(k)}([0, \infty) \times \mathbb{R})$$

$$\Delta \mathbf{u} = \begin{cases} \Delta_k u_r^k - \frac{1}{r^2} u_r^k - \frac{2ik}{r} u_\theta^k & \hat{\mathbf{r}} \\ \Delta_k u_\theta^k - \frac{1}{r^2} u_\theta^k + \frac{2ik}{r} u_r^k & \hat{\boldsymbol{\theta}} \\ \Delta_k u_z^k & \hat{\mathbf{z}} \end{cases}$$

Next Steps

Persistence result for Galerkin approximation

Fully localised patterns from Turing-Hopf bifurcations

Rigorous verification of envelope equations in [HL24]

Stability of localised planar structures (spots, patterns)

Planar invasion behaviour (spreading of pattern patches)

Localised 3D patterns - spherical harmonics expansions

$$u(r, \theta, \varphi) = \sum_{n=0}^{\infty} \sum_{m=-n}^n u_{n,m}(r) Y_n^m(\theta, \varphi)$$

3D function spaces?